

## Second Quantization

Reading: Condensed Matter Field Theory, Altland and Simons (2006)

### 1. Commutator algebra

Show that for any three operators  $A, B, C$ , the following relations hold (where  $[A, B] = AB - BA$  and  $\{A, B\} = AB + BA$ ):

$$[AB, C] = A[B, C] + [A, C]B = A\{B, C\} - \{A, C\}B \quad (1)$$

$$[A, BC] = [A, B]C + B[A, C] = \{A, B\}C - B\{A, C\} \quad (2)$$

### 2. Eigenstates of Diagonal Hamiltonian

Consider the diagonal Hamiltonian  $H = \sum_k \varepsilon_k b_k^\dagger b_k$ , where  $k$  is a discrete index,  $\varepsilon_k$  is a corresponding discrete energy, and  $b_k$  and  $b_k^\dagger$  are creation and annihilation operators satisfying either bosonic or fermionic canonical commutation relations:

$$\text{bosonic : } [b_k, b_{k'}^\dagger] = \delta_{kk'}, \quad [b_k, b_{k'}] = 0, \quad [b_k^\dagger, b_{k'}^\dagger] = 0, \quad (3)$$

$$\text{fermionic : } \{b_k, b_{k'}^\dagger\} = \delta_{kk'}, \quad \{b_k, b_{k'}\} = 0, \quad \{b_k^\dagger, b_{k'}^\dagger\} = 0. \quad (4)$$

- (i) Show that  $[H, b_k^\dagger] = \varepsilon_k b_k^\dagger$ .
- (ii) Show that if  $|E\rangle$  is an eigenstate of  $H$  with eigenenergy  $E$ , then  $b_k^\dagger|E\rangle$  is also an eigenstate, with energy  $E + \varepsilon_k$ .
- (iii) Let  $|0\rangle$  be the “vacuum state”, defined by the condition that  $b_k|0\rangle = 0$  for all  $k$ . Write down a general expression for the eigenstates of  $H$ , expressed in terms of products of creation operators acting on  $|0\rangle$ . What is the corresponding eigenenergy?

### 3. Tight-binding chain

The Hamiltonian for a periodic tight-binding chain of length  $L$  is given by

$$H_{\text{chain}} = -t \sum_{n=1}^L \left( a_n^\dagger a_{n+1} + a_{n+1}^\dagger a_n \right) \quad (5)$$

where  $t$  is the hopping matrix element  $t$  between neighboring sites  $n$  and  $n+1$ , and  $a_n^\dagger$  creates a fermion on site  $n$ , and the set of operators  $\{a_n^\dagger, a_n; n = 1, \dots, L\}$  satisfies canonical anticommutation relations of the form of Eq. (4). We assume periodic boundary conditions, i.e. we make the identification  $a_{L+n}^\dagger \equiv a_n^\dagger$ .

The purpose of this exercise is to show that this Hamiltonian can be diagonalized by a linear transformation having the form of a discrete Fourier transformation:

$$b_k^\dagger = \frac{1}{\sqrt{L}} \sum_{n=1}^L e^{ikn} a_n^\dagger \quad (6)$$

- (i) Let us require that  $b_k^\dagger$  remains invariant under an arbitrary shift of the summation index,  $n \rightarrow n + n'$  (“translational invariance”). Show that this implies that the index  $k$  is quantized, and determine the set of allowed  $k$ -values. How many independent  $b_k^\dagger$  operators are there?
- (ii) Check that the set of  $b_k$  and  $b_k^\dagger$  operators satisfies canonical anticommutation relations [Eq. (4)]. Hint: use the identity  $\sum_{m=1}^L e^{i2\pi m/L} = 0$ .
- (iii) Show that the inverse of the transformation (6) has the form

$$a_n^\dagger = \frac{1}{\sqrt{L}} \sum_k e^{-ikn} b_k^\dagger, \quad (7)$$

where the sum is over the set of allowed  $k$ -values determined in (i).

- (iv) Show that  $b_k^\dagger$  actually is a creation operator for a 1-particle *eigenstate* of  $H$ , by showing that its commutator with the Hamiltonian  $H_{\text{chain}}$  has the form  $[H_{\text{chain}}, b_k^\dagger] = \varepsilon_k b_k^\dagger$ . Give an explicit expression for the corresponding eigenenergy  $\varepsilon_k$ .
- (v) The result of (iv) implies that the Hamiltonian can be written in the form  $H = \sum_k \varepsilon_k b_k^\dagger b_k$ . Verify this explicitly, by inserting Eq. (7) into Eq. (5) for  $H_{\text{chain}}$ , and simplifying!
- (vi) Give a formula for the ground state of a half-filled chain (total particle number =  $L/2$ ).

#### 4. Fermionic representation of spin operators

Let  $c_{n\mu}^\dagger$  be the creation operators for a set of spinful fermions labeled by a discrete index  $n$  (for sites on a chain) and a spin index  $\mu = +1$  or  $-1$ . The total spin of these fermions is described by the set of three spin operators

$$S_i \equiv \frac{\hbar}{2} \sum_{n\mu} c_{n\mu}^\dagger \sigma_{\mu\mu'}^i c_{n\mu'} \quad (i = \{x, y, z\})$$

where  $\sigma^i$  are the Pauli matrices

$$\sigma^x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \sigma^y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \quad \sigma^z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}.$$

that obey  $[\sigma_i, \sigma_j] = i2 \sum_k \varepsilon_{ijk} \sigma_k$ .

Show that the operators  $S_i$  fulfill the standard spin commutator relations

$$[S_i, S_j] = \hbar i \sum_k \varepsilon_{ijk} S_k.$$