

1. Harmonisches Potential

(1a)

Potential:  $V(x, y) = \frac{m\omega^2}{2} (x^2 + y^2)$  (1)

Zylinderkoordinaten:  $x = \rho \cos \phi, \quad y = \rho \sin \phi,$

Lagrange-Funktion:  $L = T - V$

$L = \frac{1}{2} m (\dot{\rho}^2 + \rho^2 \dot{\phi}^2) - \frac{m\omega^2}{2} \rho^2$  (2)

(1b)

$\frac{d}{dt} \frac{\partial L}{\partial \dot{\phi}} = \frac{\partial L}{\partial \phi}$

1.2

$m \frac{d}{dt} (\rho^2 \dot{\phi}) = 0$

$\Rightarrow$  Erhaltener Drehimpuls:  $L = m \rho^2 \dot{\phi}$  (3)

$\dot{\phi} = \frac{L}{m \rho^2}$  (4)

da  $L \neq L(t)$ :

Energieerhaltung:  $\int \frac{L^2}{m^2 \rho^4}$  (4)

(1c)

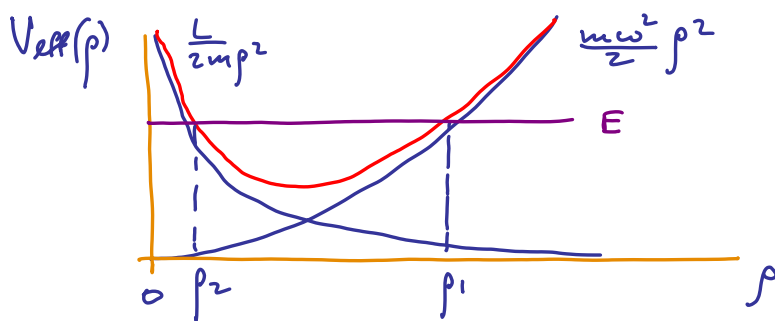
$E = \frac{1}{2} m (\dot{\rho}^2 + \rho^2 \dot{\phi}^2) + \frac{m\omega^2}{2} \rho^2$  (5a)

$= \frac{1}{2} m \dot{\rho}^2 + \underbrace{\frac{L^2}{2m \rho^2}}_{V_{eff}(\rho)} + \frac{m\omega^2}{2} \rho^2$  (5b)

mit

$$V_{\text{eff}}(\rho) = \frac{L^2}{2m\rho^2} + \frac{m\omega^2}{2}\rho^2 \quad (6)$$

1.3



(1d)

Umkehrpunkte  $\rho = \rho_1$  oder  $\rho_2$ : wenn  $\dot{\rho} = 0$ , d.h.

$$E = V_{\text{eff}}(\rho) = \frac{L^2}{2m\rho^2} + \frac{1}{2}m\omega^2\rho^2 \quad (7)$$

Umformung:  
 $\frac{2\rho^2}{m\omega^2} \times (10)$

$$\rho^4 - \frac{2E}{m\omega^2}\rho^2 + \frac{L^2}{m^2\omega^2} = 0$$

Zwei Lösungen:

$$\begin{aligned} \left\{ \begin{array}{l} \rho_1^2 \\ \rho_2^2 \end{array} \right\} &= \frac{E}{m\omega^2} \pm \sqrt{\frac{E^2}{m^2\omega^4} - \frac{L^2}{m^2\omega^2}} \\ &= \frac{E}{m\omega^2} \left[ 1 \pm \sqrt{1 - \frac{L^2\omega^2}{E^2}} \right] \end{aligned} \quad (8)$$

1.5

(1e) Für Kreisbahnen ist  $\rho_1 = \rho_2 = \sqrt{\frac{E}{m\omega^2}} \equiv \rho_0$  (9)

$$\stackrel{(8)}{\Rightarrow} \sqrt{1 - \frac{L^2\omega^2}{E^2}} = 0 \Rightarrow \boxed{E/L = \omega} \quad (10)$$

Aber:

$$L \stackrel{(3)}{=} m\rho_0^2 \dot{\phi} \quad (11)$$

$$\Rightarrow \boxed{\dot{\phi}} \stackrel{(11)}{=} \frac{L}{m\rho_0^2} \stackrel{(9)}{=} \frac{L}{m} \left( \frac{m\omega^2}{E} \right) \stackrel{(10)}{=} \boxed{\omega} \quad (12)$$

(1f) Bonus: Radiale Kraft  $F_p = -\frac{\partial V}{\partial \rho} = -m\omega^2 \rho_0$

1.5

Zentripetalkraft  $F_c = \frac{mv^2}{\rho_0} = m\rho_0 \dot{\phi}^2$   
 denn:  $v = \rho_0 \dot{\phi}$

$|F_p| = |F_c| \Rightarrow m\rho_0 \dot{\phi}^2 = m\omega^2 \rho_0$

$\Rightarrow \dot{\phi} = \omega$  (3)

Eingesetzt in: (5a), mit  $\dot{\rho} = 0$   
 $E = \frac{1}{2} m v^2 + \frac{1}{2} m \omega^2 \rho_0^2 = m \omega^2 \rho_0^2$   
 $v = \rho_0 \dot{\phi}$

(3)  $L = m \rho_0^2 \dot{\phi} \stackrel{(13)}{=} m \rho_0^2 \omega$

$\Rightarrow E/L = \omega \checkmark$  konsistent mit (4)

2. Kanonische Transformation:

2.1

Sei  $\omega, \alpha = \text{konst.}$

$H(q, p) = \frac{p^2}{2} + \frac{\omega^2 q^2}{2} + \alpha p q$  (1)

(2a)  $\dot{q} = \frac{\partial H}{\partial p} = p + \alpha q$  (2a)

$\dot{p} = -\frac{\partial H}{\partial q} = -\omega^2 q - \alpha p$  (2b)

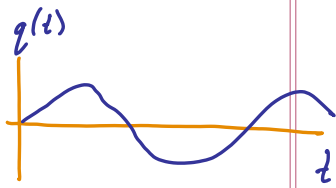
(2b)  $\frac{d}{dt}(2a):$   
 $\ddot{q} = \dot{p} + \alpha \dot{q} \stackrel{(2b)}{=} (-\omega^2 q - \alpha p) + \alpha \dot{q}$   
 $= -\omega^2 q + \alpha \dot{q} - \alpha [\dot{q} - \alpha q] \stackrel{(2a)}{=}$   
 $= -(\omega^2 - \alpha^2) q$

$\ddot{q} + \Omega^2 q = 0$  (3)

$\Omega^2 = \omega^2 - \alpha^2$  (4)

(2c) Anfangsbedingungen:  $q(0) = 0, \quad \dot{q}(0) = v_0 > 0$

2.2

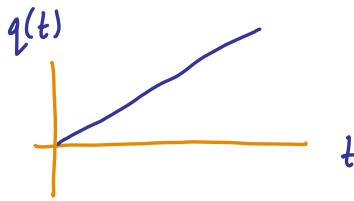


(i)  $\Omega^2 > 0$

Lösung (nicht verlangt)

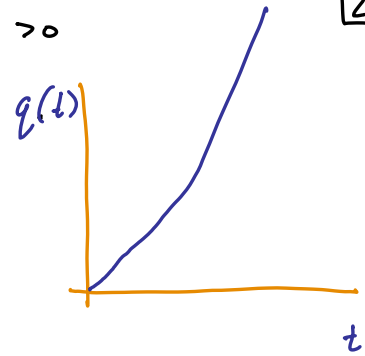
$$q(t) = \frac{v_0}{\Omega} \sin \Omega t$$

$$\Omega = \sqrt{\omega^2 - \alpha^2} > 0$$



(ii)  $\Omega^2 = 0$

$$q(t) = v_0 t$$



(iii)  $\Omega^2 < 0$

$$q(t) = \frac{v_0}{\tilde{\Omega}} \sinh \tilde{\Omega} t$$

$$\tilde{\Omega} = \sqrt{\alpha^2 - \omega^2} > 0$$

(2d) Erzeugende für kanonische Transformation:

2.3

$$F_2(q, P) = qP - \frac{A}{2} q^2 \quad (5)$$

Neue Koordinate:  $Q = \frac{\partial F_2}{\partial P} = q \quad (6a)$

Alter Impuls:  $P = \frac{\partial F_2}{\partial q} = P - Aq \quad (6b)$

Also:

$$q(Q, P) \stackrel{(6a)}{=} Q \quad (7a)$$

$$P(Q, P) \stackrel{(6b)}{=} P - AQ \quad (7b)$$

Transformierte Hamiltonfunktion:

2.4

$$K(Q, P) = H(q, p) + \frac{\partial F_2}{\partial t}$$

↖ = 0 hier

$$\stackrel{(1)}{=} \frac{P^2}{2} + \frac{\omega^2 q^2}{2} + \alpha p q$$

$$\stackrel{(10)}{=} \frac{1}{2}(P - A Q)^2 + \frac{\omega^2 Q^2}{2} + \alpha(P - A Q)Q$$

$$= \frac{1}{2} P^2 + \frac{1}{2}(\omega^2 + A^2 - \alpha A) Q^2 + P Q (-A + \alpha)$$

⇒ wähle  $A = \alpha$ , dann (8)

$$K(Q, P) = \frac{1}{2} P^2 + \frac{1}{2}(\omega^2 - \alpha^2) Q^2 \quad (9)$$

Kanonische Bewegungsgl:

$$\dot{Q} = \frac{\partial K}{\partial P} \stackrel{(9)}{=} P \quad (10a)$$

$$\dot{P} = -\frac{\partial K}{\partial Q} \stackrel{(9)}{=} -\frac{1}{2}(\omega^2 - \alpha^2) Q \quad (10b)$$

⇒

$$\ddot{Q} \stackrel{(10a)}{=} \dot{P} \stackrel{(10b)}{=} -\frac{1}{2}(\omega^2 - \alpha^2) Q$$

✓  
konsistent mit (6),  
wie erwartet,  
denn  $q = Q$ .

Check: Nicht verlangt:

sind ursprüngliche Bewegungsgleichungen für  $p, q$  erfüllt,  
falls die für  $P, Q$  erfüllt sind?

$$\dot{q} \stackrel{?}{=} P + \alpha q$$

$$\dot{Q} = (P - \alpha Q) + \alpha Q \stackrel{?}{=} P$$

$$\dot{p} = -\omega^2 q - \alpha p$$

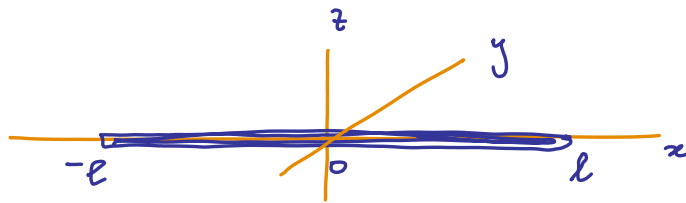
$$(\dot{p} - \alpha \dot{Q}) = -\omega^2 Q - \alpha(P - \alpha Q)$$

$$\dot{p} = -(\omega^2 - \alpha^2) Q \quad \checkmark$$

### 3. Stab-Feder System

3.1

(3a)



1-Dimensionale  
Massendichte:

$$\rho(x) = \frac{m}{2l} \Theta(l - |x|) \quad (1)$$

Drehmoment:

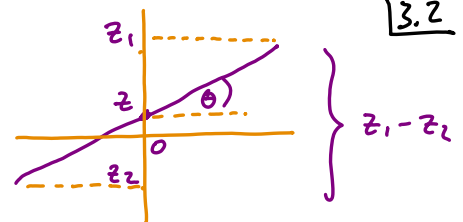
$$I = \frac{m}{2l} \int_{-l}^l dx x^2 = \frac{m}{2l} \left. \frac{1}{3} x^3 \right|_{-l}^l$$

$$= \frac{m}{2l} \frac{1}{3} 2 l^3$$

$$I = \frac{1}{3} m l^2 \quad (2)$$

(3b)  $z = \frac{z_1 + z_2}{2}, \quad \sin \theta = \frac{z_1 - z_2}{2l}$

$$\left. \begin{aligned} z_1 &= z + l \sin \theta \\ z_2 &= z - l \sin \theta \end{aligned} \right\} (3)$$



Pot. Energie:  $V(z_1, z_2) = \frac{k_1}{2} z_1^2 + \frac{k_2}{2} z_2^2$

$$= \frac{k_1}{2} (z^2 + 2z l \sin \theta + l^2 \sin^2 \theta) + \frac{k_2}{2} (z^2 - 2z l \sin \theta + l^2 \sin^2 \theta) \quad (4)$$

linearisieren:

$$\sin \theta \approx \theta$$

$$\approx \frac{1}{2} (k_1 + k_2) z^2 + (k_1 - k_2) l z \theta + \frac{1}{2} (k_1 + k_2) l^2 \theta^2 \quad (5)$$

Sei:  $K = k_1 + k_2$   
 $k = k_1 - k_2$

$$V = \frac{1}{2} K (z^2 + l^2 \theta^2) + k l z \theta \quad (6)$$

(3c)  $T = T_{sp} + T_{rot} = \frac{1}{2} m \dot{z}^2 + \frac{1}{2} I \dot{\theta}^2$  3.3  
(7)  
 $L = \frac{1}{3} m l^2$

(3d) Lagrange:  $L = T - V$

Bewegungsgleichungen:  $\frac{d}{dt} \frac{\partial L}{\partial \dot{q}} = \frac{\partial L}{\partial q}$

für  $z$ :  $m \ddot{z} = -(Kz + kl\theta)$

$\frac{1}{3} m l^2 \ddot{\theta} = -(Kl^2 \theta + klz)$

Matrix-Form:

$$\begin{pmatrix} m & 0 \\ 0 & \frac{1}{3} m \end{pmatrix} \begin{pmatrix} \ddot{z} \\ \ddot{\theta} \end{pmatrix} = - \begin{pmatrix} K & kl \\ kl/l & K \end{pmatrix} \begin{pmatrix} z \\ \theta \end{pmatrix}$$
 (8)

mit  $K = k_1 + k_2$ ,  $K = k_1 - k_2$

(3e)

Charakteristisches  
Polynom:

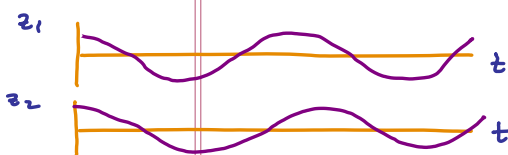
$$(K - \omega^2 m) (K - \omega^2 \frac{1}{3} m) - k^2 = 0$$
 3.4  
(9)

(3f) Sei  $k_1 = k_2 \Rightarrow k = 0$ .

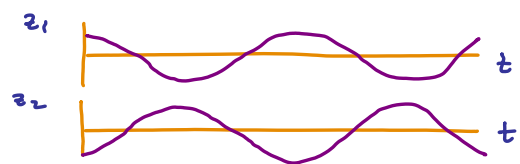
Schwerpunktschwingungen:  $\omega_1 = \sqrt{K/m}$ ,  $\begin{pmatrix} z_0 \\ \theta_0 \end{pmatrix}_1 = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$

Winkelschwingungen:  $\omega_2 = \sqrt{3K/m}$ ,  $\begin{pmatrix} z_0 \\ \theta_0 \end{pmatrix}_2 = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$

(3g)



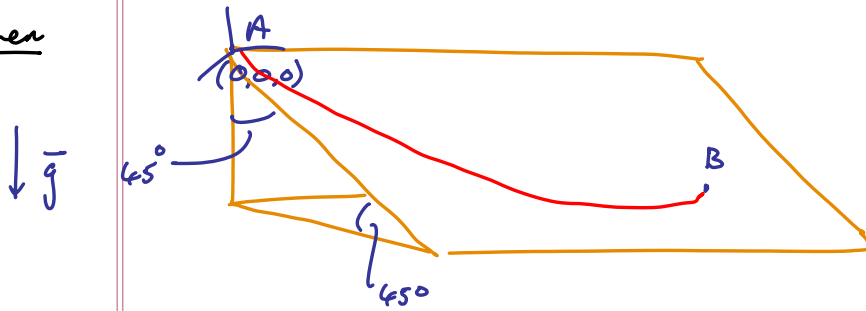
Mode 1



Mode 2

4. Ski-Rennen

(4.1)



(4a) Zwangsbedingung:  $-z = x \tan \theta = x$  (1)

(4b) Energie erhalten:  $E = \frac{1}{2} m v^2 + m g (z(t) - z_A) \equiv 0$   
 $v = \sqrt{-2gz} = \sqrt{2gx'} = v$  (2)

(4c) Fahrzeit:  $T = \int_A^B dt = \int_A^B ds \frac{1}{v}$

Aber:

$$ds = \sqrt{dx'^2 + dy^2 + dz^2}$$

(4.2)

mit  $dz = -dx$

$$= \sqrt{2dx^2 + dy^2} = dy \sqrt{1 + 2 \underbrace{\left(\frac{dx}{dy}\right)^2}_{\equiv x'^2}}$$

⇒

$$T = \int_0^{y_B} dy \frac{\sqrt{1 + 2x'^2}}{\sqrt{2gx}}$$

(3)

← keine  $y$ -Abhängigkeit

(4d) Erhaltene Größe:

$$c = F - x' \frac{\partial F}{\partial x'}$$

$$= \sqrt{\frac{1 + 2x'^2}{2gx}} - x' \frac{2x'}{(1 + 2x'^2) 2gx}$$

$$= \frac{1}{\sqrt{(1 + 2x'^2) 2gx}}$$



4.3

$$2x'^2 + 1 = \frac{1}{2gx c^2}$$

$$\frac{dx}{dy} = x' = \pm \sqrt{\frac{1}{4gc^2 x} - \frac{1}{2}} \quad (4)$$

$$y = \int dy = \pm \int_0^x dx \sqrt{\frac{4gc^2 x}{1 - 2gc^2 x}} + a$$

(4e) Substitution:  $x = \frac{1}{2c^2 g} \sin^2 \varphi/2 = \frac{1}{4c^2 g} (1 - \cos \varphi)$

$$dx = \frac{1}{2c^2 g} \sin \varphi \cos \varphi/2 d\varphi$$

$$y = \pm \int_0^{\varphi} \frac{d\varphi}{2c^2 g} \sin \varphi \cos \varphi/2 \frac{\sqrt{2} \sin \varphi/2}{\sqrt{1 - \sin^2 \varphi/2}} + a$$

$$= \pm \frac{1}{\sqrt{2} c^2 g} \int d\varphi \underbrace{\sin^2 \varphi/2}_{\frac{1}{2}(\varphi/2 - \cos \varphi/2 \sin \varphi/2)} + a \quad (4.4)$$

$\frac{1}{2} \sin \tilde{\varphi}$

$$y = \pm \left( \frac{1}{\sqrt{2} c^2 g} \right) \frac{1}{4} (\varphi - \sin \varphi) + a \quad (5)$$

$x = 0 \Rightarrow \varphi = 0$  und  $y = 0 \Rightarrow a = 0.$

(4f)

Also gilt:  $x = \frac{1}{4c^2 g} (1 - \cos \varphi) = r_x (1 - \cos \varphi) = x$

$$y = \pm \left( \frac{1}{\sqrt{2} 4c^2 g} \right) (\varphi - \sin \tilde{\varphi}) = \pm r_y (\tilde{\varphi} - \cos \varphi) = y \quad (6)$$

mit  $r_x / r_y = \sqrt{2}$

Nicht verlangt:

Durch Einsetzen der Endpunktwerte  $(x_B, y_B)$ , werden  $\varphi_B$  und  $c$

festgelegt:

$$x_B = \frac{1}{4c^2g} (1 - \cos\varphi_B)$$

$$y_B = \frac{1}{\sqrt{2} 4c^2g} (\varphi_B - \sin\varphi_B)$$