

2. Durchhängende Hochspannungsleitung

$$a) \quad U[y(x)] = \int_{-x_0}^{x_0} \frac{m}{\ell} g y(x) \sqrt{1 + [y'(x)]^2} dx$$

$$b) \quad L[y(x)] = \int_{-x_0}^{x_0} \sqrt{1 + [y'(x)]^2} dx = \ell$$

$$c) \quad \begin{cases} U[y(x)] = \text{minimal} \\ L[y(x)] = \ell \end{cases}$$

$$U[y(x)] + \lambda (L[y(x)] - \ell) = \text{minimal}$$

$$\int_{-x_0}^{x_0} \left\{ \frac{mg}{\ell} y(x) \sqrt{1 + [y'(x)]^2} + \lambda \left(\sqrt{1 + [y'(x)]^2} - \frac{\ell}{2x_0} \right) \right\} dx = \text{minimal}$$

$$\delta \int_{-x_0}^{x_0} F(y, y', x) dx = 0$$

$$F(y, y', x) = \left(\frac{mg}{\ell} y + \lambda \right) \sqrt{1 + y'^2} - \lambda \frac{\ell}{2x_0}$$

$$d) \quad F - y' \frac{\partial F}{\partial y'} = C$$

$$\frac{\partial F}{\partial y'} = \left(\frac{mg}{\ell} y + \lambda \right) \frac{y'}{\sqrt{1 + y'^2}}$$

$$F - y' \frac{\partial F}{\partial y'} = \left(\frac{mg}{l} y + \lambda \right) \sqrt{1 + y'^2} - \lambda \frac{l}{2x_0} -$$

$$- y' \left(\frac{mg}{l} y + \lambda \right) \frac{y'}{\sqrt{1 + y'^2}} = C$$

$$\left(\frac{mg}{l} y + \lambda \right) \frac{1}{\sqrt{1 + y'^2}} - \lambda \frac{l}{2x_0} = C$$

$$y' = \pm \sqrt{\left(\frac{\frac{mg}{l} y + \lambda}{C + \frac{\lambda l}{2x_0}} \right)^2 - 1}$$

e)

$$\frac{dy}{dx} = \pm \sqrt{\left(\frac{\frac{mg}{l} y + \lambda}{C + \frac{\lambda l}{2x_0}} \right)^2 - 1}$$

$$\int \frac{dy}{\sqrt{\left(\frac{\frac{mg}{l} y + \lambda}{C + \frac{\lambda l}{2x_0}} \right)^2 - 1}} = \pm x + \alpha$$

$$\frac{C + \frac{\lambda l}{2x_0}}{\frac{mg}{l}} \int \frac{dz}{\sqrt{z^2 - 1}} = \pm x + \alpha, \quad z = \frac{\frac{mg}{l} y + \lambda}{C + \frac{\lambda l}{2x_0}}$$

$$\operatorname{arccosh} \left(\frac{\frac{mg}{l} y + \lambda}{C + \frac{\lambda l}{2x_0}} \right) = \frac{\frac{mg}{l} (\alpha \pm x)}{C + \frac{\lambda l}{2x_0}}$$

$$\frac{mg}{l} y + \lambda = \left(C + \frac{\lambda l}{2x_0} \right) \cosh \left(\frac{\frac{mg}{l} (\alpha \pm x)}{C + \frac{\lambda l}{2x_0}} \right)$$

$$y = \frac{l}{mg} \left\{ \left(c + \frac{\lambda l}{2x_0} \right) \cosh \left(\frac{\frac{mg}{l} (a \pm x)}{c + \frac{\lambda l}{2x_0}} \right) - \lambda \right\} \quad (3)$$

One could also absorb $\frac{\lambda l}{2x_0}$ into the c -constant. An equivalent solution is, therefore,

$$y = \frac{l}{mg} \left\{ C \cosh \left(\frac{\frac{mg}{l} (a \pm x)}{C} \right) - \lambda \right\}$$

Furthermore, $\frac{l}{mg} C$ could be also taken into C .

$$y = C \cosh \left(\frac{a \pm x}{C} \right) - \frac{\lambda l}{mg}$$

All three solutions are equivalent from point of view of solving differential equations.

$$f) \quad y(-x_0) = 0$$

$$y(x_0) = 0$$

$$\int_{-x_0}^{x_0} \sqrt{1 + [y'(x)]^2} dx = l$$

Facultative: Let's apply these conditions.

$$0 = C \cosh \left(\frac{a \pm x_0}{C} \right) - \frac{\lambda l}{mg}$$

$$0 = C \cosh \left(\frac{a \mp x_0}{C} \right) - \frac{\lambda l}{mg}$$

$$\int_{-x_0}^{x_0} \sqrt{1 + \sinh^2 \left(\frac{a \pm x}{C} \right)} dx = l$$

From the first two it follows that $a = 0$. Then, we are left with:

$$0 = C \cosh \left(\frac{x_0}{C} \right) - \frac{\lambda l}{mg}$$

$$\int_{-x_0}^{x_0} \cosh\left(\frac{x}{c}\right) dx = l$$

(4)

The first equation expresses λ via c

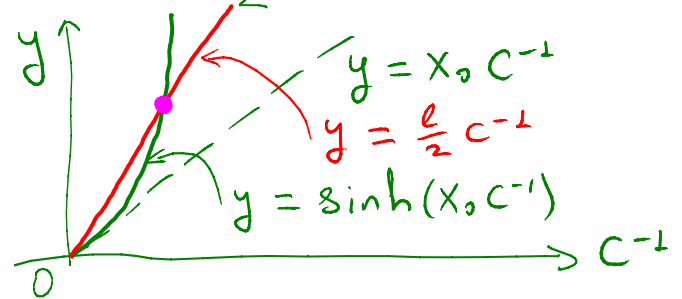
$$\lambda = \frac{mg}{\ell} c \cosh\left(\frac{x_0}{c}\right)$$

The second equation is for determining c

$$2c \sinh\left(\frac{x_0}{c}\right) = l$$

This equation is transcendental, but it has a finite solution for c , provided $x_0 < \frac{l}{2}$:

$$\sinh\left(\frac{x_0}{c}\right) = \frac{l}{2c}$$



$$\begin{cases} y(x) = c \cosh\left(\frac{x}{c}\right) - c \cosh\frac{x_0}{c} \\ \sinh\left(\frac{x_0}{c}\right) = \frac{l}{2c} \end{cases}$$

3. Überkippen eines Würfels

$$\begin{aligned} a) \quad I_{xx}^{sp} &= \int_{-\frac{a}{2}}^{\frac{a}{2}} dx \int_{-\frac{a}{2}}^{\frac{a}{2}} dy \int_{-\frac{a}{2}}^{\frac{a}{2}} dz \frac{m}{a^3} (y^2 + z^2) = \\ &= \frac{m}{a^3} \left(a^2 \frac{1}{3} y^3 \Big|_{-\frac{a}{2}}^{\frac{a}{2}} + a^2 \frac{1}{3} z^3 \Big|_{-\frac{a}{2}}^{\frac{a}{2}} \right) = \frac{m a^2}{6} \end{aligned}$$

$$I_{xy}^{sp} = - \int_{-\frac{a}{2}}^{\frac{a}{2}} dx \int_{-\frac{a}{2}}^{\frac{a}{2}} dy \int_{-\frac{a}{2}}^{\frac{a}{2}} dz \frac{m}{a^3} xy = 0$$

$$I_{ij}^{sp} = \begin{pmatrix} \frac{m a^2}{6} & 0 & 0 \\ 0 & \frac{m a^2}{6} & 0 \\ 0 & 0 & \frac{m a^2}{6} \end{pmatrix},$$

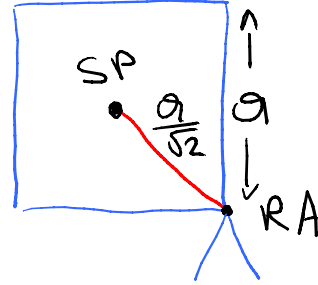
$$I_{xx}^{sp} = I_{yy}^{sp} = I_{zz}^{sp} = \frac{m a^2}{6}$$

b) All main moments of inertia are equal to each other. Therefore, one can choose the direction of one main axis arbitrarily. The tensor is diagonal in any coordinate frame.

c)
$$I_{ij} = I_{ij}^{SP} + m(\underline{r}^2 \delta_{ij} - r_i r_j)$$

$$\underline{r} = \left(\frac{a}{2}, \frac{a}{2}, 0 \right)$$

$$I_{ij} = \begin{pmatrix} \frac{5}{12} m a^2 & -\frac{m a^2}{4} & 0 \\ -\frac{m a^2}{4} & \frac{5}{12} m a^2 & 0 \\ 0 & 0 & \frac{2}{3} m a^2 \end{pmatrix}$$



Alternatively,

$$I_{zz} = I_{zz}^{SP} + m \left(\frac{a}{\sqrt{2}} \right)^2 = \frac{m a^2}{6} + \frac{m a^2}{2} = \frac{2}{3} m a^2$$

d)
$$L_{\text{before}} = m v_0 \frac{a}{2}$$

$$L_{\text{after}} = \frac{2}{3} m a^2 \omega_1$$

$$L_{\text{before}} = L_{\text{after}}$$

$$\cancel{m} v_0 \frac{\cancel{a}}{2} = \frac{2}{3} \cancel{m} a^2 \omega_1$$

$$\omega_1 = \frac{3}{4} \frac{v_0}{a}$$

e)
$$v_1 = \omega_1 \frac{a}{\sqrt{2}} = \frac{3}{4\sqrt{2}} v_0$$

$$E_k^{\text{before}} = \frac{m}{2} v_0^2$$

$$E_{\kappa}^{\text{after}} = \frac{1}{2} I_{zz} \omega_1^2 = \frac{1}{2} \frac{2}{3} m a^2 \left(\frac{3}{4} \frac{v_0}{a} \right)^2 = \frac{3}{16} m v_0^2 \quad (6)$$

$$E_{\kappa}^{\text{before}} - E_{\kappa}^{\text{after}} = \frac{m}{2} v_0^2 - \frac{3}{16} m v_0^2 = \frac{5}{16} m v_0^2$$

$$f) \quad \frac{1}{2} I_{zz} \omega_1^2 = m g \left(\frac{a}{\sqrt{2}} - \frac{a}{2} \right)$$

$$\frac{1}{2} \frac{2}{3} m a^2 \frac{9}{16} \frac{v_0^2}{a^2} = m g a \frac{\sqrt{2}-1}{2}$$

$$v_0^2 = \frac{8}{3} (\sqrt{2}-1) g a$$

$$v_0 = v_g = \sqrt{\frac{8(\sqrt{2}-1)}{3} g a}$$